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Ionospheric Vertical Correlation Distances: Estimation From ISR Data, Analysis, and Implications For Ionospheric Data Assimilation

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Abstract The construction of the background covariance matrix is an important component of ionospheric data assimilation algorithms, such as Ionospheric Data Assimilation Four-Dimensional (IDA4D). It is a matrix that describes the correlations between all the grid points in the model domain and determines the transition from the data-driven to model-driven regions. The vertical component of this matrix also controls the shape of the assimilated electron density profile. To construct the background covariance matrix, the information about the spatial ionospheric correlations is required. This paper focuses on the vertical component of the model covariance matrix. Data from five different incoherent scatter radars (ISR) are analyzed to derive the vertical correlation lengths for the International Reference Ionosphere (IRI) 2016 model errors, because it is the background model for IDA4D. The vertical distribution of the correlations is found to be asymmetric about the reference altitude around which the correlations are calculated, with significant differences between the correlation lengths above and below the reference altitude. It is found that the correlation distances not only increase exponentially with height but also have an additional bump-on-tail feature. The location and the magnitude of this bump are different for different radars. Solar flux binning introduces more pronounced changes in the correlation distances in comparison to magnetic local time (MLT) and seasonal binning of the data. The latitudinal distribution of vertical correlation lengths is presented and can be applied to the construction of the vertical component of the background model covariance matrix in data assimilation models that use IRI or similar empirical models as the background.

1. Introduction

The background error covariance matrix \( \hat{P} \) is an important component of ionospheric data assimilation. It describes the variance of the background model used for the assimilation and how the errors of this model correlate between any two grid points. Each element of the model error covariance matrix \( \hat{P} \) can be expressed in the following way:

\[
\hat{P}_{ij} = \sigma_i \sigma_j r_{ij},
\]

where \( \sigma_i \) and \( \sigma_j \) are the standard deviations of the forecast model values at \( i \)th and \( j \)th grid points and \( r_{ij} \) is the linear correlation coefficient between these errors.

In practice, it is close to impossible to find the real representation of this matrix for several reasons. First, the true values of the electron density at all grid points on the globe are unknown. Second, the error covariance matrix \( \hat{P} \) is usually a very large matrix of size \( n \times n \), where \( n \) is the number of grid points in the assimilation. Even if, hypothetically, the truth would be known, it can be too computationally demanding to calculate \( \hat{P} \) for each time step of the assimilation. Therefore, in practice, this error covariance matrix needs to be modeled a priori. Typically, three assumptions are used for modeling the error covariance matrix (Aa et al., 2015, 2016; Bust & Crowley, 2007; Bust & Datta-Barua, 2014; Bust et al., 2001, 2004; Coker et al., 2001; Yue, Wan, Liu, Zheng, et al., 2007; Yue et al., 2011). First, the spatial correlation is assumed to be separable horizontally and vertically. Second, the vertical correlations are represented by a Gaussian. Third, the horizontal correlations are modeled by an elliptical Gaussian in geomagnetic coordinates.

Using these three assumptions, the construction of the error covariance matrix \( \hat{P} \) can be separated into the construction of three matrices: the model variance matrix \( \hat{V} \), the horizontal correlation matrix \( \hat{C}_{h\text{or}} \), and the
vertical correlation matrix $\tilde{C}_{\text{ver}}$. The elements $\tilde{P}_{ij}$ can then be found using element by element multiplication:

$$\tilde{P} = \tilde{V} \circ \tilde{C}_{\text{hor}} \circ \tilde{C}_{\text{ver}}.$$  \hspace{1cm} (2)

Forsythe et al. (2020) discussed in detail the construction of the horizontal correlation matrix $\tilde{C}_{\text{hor}}$. The construction of the vertical correlation matrix $\tilde{C}_{\text{ver}}$ is treated in this paper.

Bust et al. (2001) first suggested to model the vertical correlation matrix as a Gaussian

$$\tilde{C}_{ij} = \exp \left[ -\frac{(z_i - z_j)^2}{L_z^2} \right].$$  \hspace{1cm} (3)

where $z$ is the height and $L_z$ is the vertical correlation length for grid points $i$ and $j$.

Previously, the vertical correlation length $L_z$ was approximated by the ionospheric scale height (Bust et al., 2004) and has not been derived for any particular background model. In practice, and as will be shown in this paper, the height variation of vertical correlation distances derived from model errors can be more complicated than the ionospheric scale height variation.

In this study, the parameters that describe the latitudinal distribution of the vertical ionospheric correlations are derived for the first time. Importantly, this is the first study that is dedicated to the vertical correlation of the International Reference Ionosphere 2016 (IRI) (Bilitza et al., 2017) model errors using multiple instruments located at different latitudinal regions. The results of this study can be directly applied to the construction of the error covariance matrix for various data assimilation models (Forsythe, 2020) that use IRI or other similar empirical models as the background.

2. Experimental Setup

The data from five incoherent scatter radars (ISR) were analyzed in this study to calculate the vertical correlation lengths. Figure 1a shows the locations of Jicamarca (JRO), Arecibo (ARO), Millstone Hill (MLH), Poker Flat ISR (PFISR), and Resolute Bay North ISR (RISR-N). The yellow lines show the 80° N, 60° N, 30° N, 30° S, and 60° S geomagnetic latitudes, based on the altitude-adjusted corrected geomagnetic (AACGM) coordinate system with 2010 coefficients (Shepherd, 2014). In this study we use all available data where the radars observed the vertical distribution of the electron density for the time period starting from the year of 2000. The following ISR modes were chosen: Long Pulse (LP) Mode for RISR-N and PFISR, Oblique Mode Faraday Rotation With Uncoded LP (Hybrid 2) for JRO, and Coded LP Mode mode for ARO. For the MLH
**Figure 2.** Observed EDP from the JRO radar on 8 September 2010 at 20:05–20:10 UT shown in black solid lines. The error bars are also shown but are usually significantly smaller than the circles representing the data points. The IRI profile for this time period is shown in yellow. The dark gray area shows the altitudinal range where IRI profile is greater than half of IRI $N_mF_2$ (shown with the yellow circle). The $N_mF_2$ of the observed profile is found as the maximum density within the gray area and is shown with orange circle. Orange and cyan lines show the top fit and bottom fit with scale heights $H_m$ and correlation coefficients to the data points $R$ also shown in the panel.

ISR, we use the data provided in the “gridded data filtered to a uniform spatial and temporal grid” format, which are derived from raw zenith measurements.

For all radars listed above, only the vertical profiles were selected for the analysis presented here. For PFISR the vertically oriented beam was tagged as 64280. For RISR-N the beam with elevation angle of 75° and azimuth of 26° was selected over the beam with 90° elevation because the vertical beam has a grating lobe issue. The data from MLH and ARO used in this study were collected at elevation angles of 89.74° and 90°, respectively. The lowest elevation angle for the JRO radar data was 87.06°. Figure 1b shows the data coverage for all radars. PFISR had the highest number of observations, and MLH had the most continuous coverage between 2000 and 2020. In this study, the binning of data in magnetic local time (MLT), season, and F10.7 will be only applied to MLH radar because of the continuity of the data set. All radars had different maximum ranges. The maximum ranges for JRO, ARO, MLH, PFISR, and RISR-N are 1,635, 687, 547.85, 673.3, and 692.34 km, respectively.

### 3. Data Preprocessing

The first step of the analysis included the uniform gridding of the ISR data into 5-min intervals, excluding MLH that was already gridded into 15-min intervals. Next, the two-layer Chapman function (Rishbeth & Garriott, 1969) was fitted to each of the individual profiles. Figure 2 shows one example of the fitted Chapman model to the JRO measurements on 8 September 2010 at 20:05–20:10 UT. The measured electron density profile (EDP) with uncertainties is shown in black. In this particular time frame the uncertainties are very small and are visible only near the peak of the profile. The IRI EDP for the location of the radar and for the same time period as the measurement is shown as the yellow curve in Figure 2. The altitudinal range where the IRI density is greater than half of $N_mF_2$ (shown with the yellow circle) was determined and highlighted in the figure with dark gray color background. In case the calculated peak of the profile was located at the upper (lower) boundary of the gray area, the altitudinal range was shifted up (down) by 100 km (50 km) up (down) and the location of the density maximum was found again. This method was applied to avoid cases where the density maximum is located in the $E$ region due to sporadic $E$ events and particle precipitation. Once the $N_mF_2$ value was determined, the Chapman model described by Equation 4

$$N_e(h) = N_mF_2 \exp \left[ \frac{0.5 \left( 1 - \frac{h - h_mF_2}{H(h)} \right)}{\exp \left( \frac{-h - h_mF_2}{H(h)} \right)} \right]$$

was fitted to the data using the least squares method by Markwardt (2009). For the topside ionosphere, the scale height is given by:

$$H(h) = A_1(h - h_mF_2) + H_{m1}.$$  \hspace{1cm} (5)

This approach to modeling the topside has been applied before, demonstrating acceptable performance below 1,200-km altitudes (dos Santos Prol et al., 2019). For the bottomside ionosphere, the expression for the scale height is as follows:

$$H(h) = A_2(h - h_mF_2) + H_{m2}.$$  \hspace{1cm} (6)

Here, $h$ is the height and $A_1, A_2, H_{m1}$, and $H_{m2}$ are the fitting coefficients. Additionally, we also constrained the $A_1$ parameter to positive values only. The orange and cyan lines in Figure 2 show the fitted profile for the topside and bottomside, respectively. Since this study focuses only on the $F$ region portion of the profile, the omission of the $E$ region layer in the model is deliberate and does not impact our results. The following criteria were used to filter out poor quality profiles: The number of data points for the topside is below 5
Figure 3. (a–e) Distributions of IRI model errors normalized by the total amount of EDP profiles for each radar.

and for the bottomside is below 3, the $H_{m1}$ or the $H_{m2}$ is less then 0 or greater than 150, and the correlation coefficient between the data points and corresponding fitted points is less then 0.8. Additionally, several profiles that had negative scale height in the bottomside were excluded. Figure 1b showed the number of profiles that remained after the filtering procedure. Uncertainties for the data points were considered in the fitting procedure using normal weighting. In case the data point did not have any information about the error, it was set to 20% of the observed value.

Finally, the differences between the IRI profiles (obtained with default IRI model options) and the fitted profiles were calculated and are referred to as model errors hereafter. Figure 3 shows the distribution of model errors for all radars, normalized by the maximum number of EDPs for each radar. The peaks of the distributions of model errors for most altitudes are centered around zero. There exist some shifts of the distributions for the altitude range between 200 and 400 km. For example, the MLH radar $F$ region densities are lower in comparison to IRI, as shown in Figure 3c. The data-model differences are largest at JRO, which show negative model errors below 300-km altitude and positive errors above (Figure 3e). The tails of the distributions are very long, exceeding $1 \times 10^{12}$ m$^{-3}$ for RISR-N, PFISR, MLH, and JRO radars, indicating that a small fraction of EDPs strongly disagreed with the model. Interestingly, the RISR-N tails of the distributions are predominantly negative, as can be seen in Figure 3a. This is consistent with the previous studies (Bjoland et al., 2016; Themens et al., 2014) that concluded that the IRI model is biased toward an underestimation of the electron density in the polar cap.

4. Results

The correlations between the model errors at different heights were calculated using preprocessed data for all radars, without any binning of the data. Here we describe our analysis in more detail. For each ISR shown in Figure 3, the linear correlation coefficients are found between the model errors at a reference height with those at other heights. These correlations are calculated for a range of reference altitudes between 100 and 1,000 km. Figures 4a–4e show the calculated correlations for all radars, with the color bar shown in panel (a). The $x$ axis is the reference altitude, and the $y$ axis corresponds to the distance (positive is along the upward direction) from the reference point. The dashed white line at 200-km altitude indicates the beginning of the altitudinal range where the fitted Chapman function agrees well with the data points. Even if the data were available below this altitude, the $E$ region layer was not reflected in the fitted profile. The second white dashed line indicates the maximum altitude where the radar data were available. It is different for each radar. The maximum altitude of the JRO radar data is 1,635 km; this is why the second line is not shown. The solid (dashed) black line shows the contour of 0.7 correlation using only the points above (below) the reference point. This distance, where the correlation is equal to 0.7, is defined as the correlation length for the purpose of this study. In this paper we will keep using the same formalism and will define the correlation length as
Figure 4. Vertical correlations of IRI model errors and derived correlation length. (a–e) The color shows the correlations of IRI errors between reference altitude and all other points of the EDP, with the color bar shown in panel (a). Solid (dashed) black line shows the contour of 0.7 correlation for the points above (below) the reference point. Each row corresponds to the different radar, with the name of the radar shown at each panel. (f–j) Vertical correlation distances above and below the reference altitude are shown with solid and dashed black lines, respectively. White dashed lines in all panels indicate the regions driven by the radar data.

a distance where the linear correlation coefficient is equal to 0.7. Figures 4f–4j show the estimated vertical correlation length based on our formalism for each radar site. Again, the solid (dashed) line corresponds to the correlation length for the points above (below) the reference point.

From Figure 4 it is evident that the correlation lengths derived for upward and downward direction are very different and that this asymmetry needs to be taken into account for the modeling of the vertical covariance.
matrix. The correlation length derived from model errors exhibits a complex structure, with a well-defined bump-on-tail like shape, where the bump is centered at different reference altitudes for different radars.

Figure 5 compares the calculated vertical correlation lengths from different radars. The colors of the lines correspond to different radars. All correlation distances, except those computed from the JRO radar data, exhibit exponential increases starting around 200-km altitude. For the JRO radar site, the vertical correlation distances show exponential growths starting near 350 km for positive distances and 400 km for negative distances from the reference altitude, as shown in Figure 5. RISR-N correlation distances also start the exponential increase from 400 km for negative distances. In general, the sharpness of the exponential increase, the height of the bump, and location with respect to the reference altitude are different for each radar. Only PFISR and MLH correlation distances exhibit similar behavior.

Next, the variations of the vertical correlation distances for different diurnal, seasonal, and solar conditions are evaluated. Only the MLH data were used for this evaluation, because it was the only site that provided continuity of data coverage needed to examine the seasonal and solar flux influence on the correlation lengths. To evaluate the diurnal variation of the vertical correlation distances, the MLH data were divided into 3-hour MLT bins. In Figure 6 the colors of the lines represent different MLT bins, and the thick black line shows the correlation distance without binning. The trends look the same for all MLT bins, with minor differences in the position of the bump and the slope of the exponential increase. Both positive and negative distances from the reference altitude, shown in Figure 6, show similar behaviors with MLT.
Figure 7. Vertical correlation distances derived from IRI model errors (a) above and (b) below the reference altitude for MLH radar data binned in F10.7, indicated by the color. Thick black line shows correlation distances for data not divided into F10.7 bins.

Figure 7 shows the vertical correlation distances computed using the MLH radar data for low, moderate, and high values of F10.7 solar flux index, represented by different colors. The solar flux influence is more pronounced than the MLT variation. For the correlation distances above the reference point (Figure 7a) the high solar flux reduces the height of the bump and increases the exponential slope, whereas during low and moderate F10.7 values the bumps have a similar shape and height. The reference altitude position of the bump does not change across the different solar flux bins. For the correlation distances below the reference point (Figure 7b) the position of the bump shows similar change with solar flux, but in addition, the reference altitude of the bump is shifting as well.

Additionally, the seasonal dependence of vertical correlation distances is investigated. Figure 8 shows the vertical correlation distances for MLH radar for winter, summer, and equinox seasons. The variations with different seasons are very minor, as demonstrated in Figure 8.

5. Application to Data Assimilation

To model the vertical correlation distance for data assimilation purposes, only the unbinned data, shown in Figure 5, were considered. When more data are available for the equatorial and polar cap regions, the binning in MLT, season, and solar flux can be incorporated into the modeling.

Figure 9 shows the modeled distribution of the vertical correlation distances as a function of magnetic latitude. Figure 9a and b corresponds to the correlation distances above (below) the reference point. The color in
Figure 9. Latitudinal distribution of modeled vertical correlation length (a) above and (b) below the reference altitude. The color indicates the correlation distance, with the color bar shown on the right. The circles show the data-derived points for the locations of the ISRs, same information as in Figure 5. These results can be download at data repository (Forsythe, 2020).

Figure 9 indicates the correlation distance, with the color bar shown on the right. The correlation distances are modeled using linear interpolation between the data-derived points, shown with circles in Figure 9. These data-derived points are the same as in Figure 5, and their magnetic latitude locations correspond to the locations of the ISRs. Northern Hemisphere values are reflected into the Southern Hemisphere in the absence of Southern Hemisphere radar data.

Then each element of the vertical correlation matrix for the construction of the background covariance matrix can be modeled as

$$
\tilde{C}_{ij}^{ver} = \begin{cases} 
\exp\left(-\frac{(z_i - z_j)^2}{L_1(z_i, \lambda_i)}\right), & \text{if } z_i < z_j \\
\exp\left(-\frac{(z_i - z_j)^2}{L_2(z_i, \lambda_i)}\right), & \text{if } z_i > z_j, \\
1, & \text{if } z_i = z_j
\end{cases}
$$

(7)

where $z$ is the height, $\lambda$ is the magnetic latitude, $L_1$ and $L_2$ are functions of altitude and magnetic latitude, and subscripts $i$ and $j$ refer to the pairs of grid points. $L_1$ and $L_2$ are presented in Figures 9a and 9b and are provided in the form of metadata.

6. Discussion

6.1. The Position of the Bump

Based on previous studies, a simple exponential increase of the correlation distances with height was anticipated. For example, Yue, Wan, Liu, and Mao (2007) analyzed vertical correlations derived from day-to-day ionospheric variability using MLH data and reported an exponential increase of correlation distance with height, even though they used a fitting scheme similar to the one in this study. The increase of the ionospheric scale height was also assumed to be exponential (Bust et al., 2004; Yue et al., 2011) without a bump-on-tail local maximum. This study shows that the correlation distances exhibit a more complex
Figure 10. Correlation distances calculated from MLH data without any removal of the corresponding IRI modeled values or day-to-day variability. Solid (dashed) lines correspond to the upward (downward) direction from the reference point. Black lines show the correlation distances derived from unmodified MLH data, whereas blue, yellow, and red lines correspond to the following modifications: \( h_m F_2 + 40 \) km, \( H_m^1 + 10 \) km, and \( H_m^2 + 10 \) km, respectively.

6.2. Comparison With Vertical Correlation Length Derived From Day-to-Day Ionospheric Variability

In the previous investigation of the horizontal correlation distances (Forsythe et al., 2020), a significant difference was found between the correlation distances derived from the IRI errors and from day-to-day total electron content (TEC) variability. It was concluded that the day-to-day variability-derived correlation length cannot be employed for the modeling of the background error covariance matrix. Similarly here, to investigate the differences between the vertical correlation lengths derived using these two approaches, the following analysis is performed. The day-to-day variability of electron density (the difference between two consecutive days) was calculated using fitted EDPs for MLH, and the correlation distances were found using the same method described in section 3. Figure 11 shows the results of the comparison. The blue (red) color corresponds to the correlation distances derived from IRI model errors (day-to-day electron density variability). The differences between the two results are very minor. This suggests that the day-to-day ionospheric variability is one of the important factors that controls the vertical distribution of model error correlations.
7. Conclusions
The IRI-2016 model errors were calculated using data from five ISRs, and the vertical correlation distances were computed from the distributions of model errors. It was found that the vertical distribution of the correlations is asymmetric and that it is important to estimate the vertical correlation distance in two directions, above and below the reference point. The correlation distances increase exponentially with height and have an additional bump-on-tail enhancement. The reference altitude and height of this bump are different for all radars. The position of the bump is controlled by the $hmF2$ and $H_m$ parameters. The changes with MLT and season for MLH radar are not significant, but the solar flux binning introduces more pronounced changes (about 100-km difference in the height of the bump for high and low solar flux). The latitudinal distribution of vertical correlation length was modeled and is available at Forsythe (2020). This distribution can be applied to the construction of vertical component of the background model covariance matrix. In a future study, the horizontal and vertical correlation lengths will be implemented in the Ionospheric Data Assimilation Four-Dimensional (IDA4D) algorithm to examine their effects on the assimilation results.

Data Availability Statement
All ISR data were obtained through Madrigal Database (http://isr.sri.com/madrigal/). The vertical correlation lengths for different latitudes and heights derived in this study are available online (https://doi.org/10.5281/zenodo.3928823).

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